## Quantumphysics 2

## Exam July 2, 2010. Tentamenhal, 9.00-12.00.

$\diamond$ Write your name and student number on each sheet you use.
$\diamond$ The exam has 5 problems. The number of points for each exercise is an indication about the time you are supposed to spend to solve it.
$\diamond$ Read the problems carefully and give complete and readable answers.
$\diamond$ No books or personal notes are allowed.

## Problem 1 (22 pts)

i) Prove that $\hat{J}^{2}=\hat{J}_{-} \hat{J}_{+}+\hat{J}_{z}^{2}+\hbar \hat{J}_{z}(5 \mathrm{pts})$

Answer: We use ladder operators to rewrite $J_{x}$ and $J_{y}$. Then, we have

$$
J_{x}^{2}+J_{y}^{2}+J_{z}^{2}=\frac{J_{+} J_{-}+J_{-} J_{+}}{2}+J_{z}^{2}
$$

Using the commutator $\left[J_{+}, J_{-}\right]=J_{+} J_{-}-J_{-} J_{+}=2 \hbar J_{z}$, we have

$$
J^{2}=\hbar J_{z}+J_{-} J_{+}+J_{z}^{2}
$$

ii) Consider a particle with spin $\mathrm{s}=1 / 2$. Derive the Pauli matrices $\hat{\sigma}_{x, y, z}$. What is the relation between these matrices and the $\hat{S}_{x, y, z}$ operators?
Hint: Work in the basis with $\hat{S}_{z}$ diagonal. Use ladder operators. (4 pts)
Answer: In this basis, it is elementary to see that

$$
S_{z}=\frac{\hbar}{2}\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right)=\frac{\hbar}{2} \sigma_{z} .
$$

Using the ladder operators, we can easily build the matrix for $S_{x}$ and $S_{y}$ and consequently derive $\sigma_{x}$ and $\sigma_{y}$.
iii) What is the action of the angular momentum operators $\hat{J}^{2}$ and $\hat{J}_{z}$ on $Y_{l}^{m}$ functions? Evaluate $\int d \Omega\left(Y_{l}^{m}\right)^{*} Y_{l^{\prime}}^{m^{\prime}}$ (2 pts)

Answer:

$$
\begin{gathered}
J^{2}\left|j, m_{j}\right\rangle=\hbar^{2} j(j+1)\left|j, m_{j}\right\rangle \quad J_{z}\left|j, m_{j}\right\rangle=m_{j} \hbar\left|j, m_{j}\right\rangle \\
\int d \Omega\left(Y_{l}^{m}\right)^{*} Y_{l^{\prime}}^{m^{\prime}}=\delta_{l l^{\prime}} \delta_{m m^{\prime}}
\end{gathered}
$$

iv) Prove that the degeneracy of the H -atom energy levels $E_{n}$ is $n^{2}$ (excluding the spin degeneracy). (2 pts)

Answer: The degeneracy $\mathcal{D}$ is

$$
\mathcal{D}=\sum_{l=0}^{n-1}(2 l+1)=n^{2} .
$$

Extra: Including spin, we have to take into account the Pauli's principle, so we can accommodate 2 electrons per energy level. Thus $\mathcal{D}=2 n^{2}$.
v) Write down the quantum mechanical angular momentum operator in a coordinate representation. (2 pts)

Answer: $\hat{L}=-i \hbar \vec{r} \times \vec{\nabla}$
vi) Evaluate the commutators $\left[\hat{L}_{x}, \hat{L}_{y}\right]$ and $\left[\hat{L}_{z}, \hat{L}_{y}\right]$. Using these outcomes, evaluate $\left[\hat{L}_{z} \hat{L}_{x}, \hat{L}_{y}\right]$. (4 pts)

Answer: The first two commutators give respectively $i \hbar L_{z}$ and $-i \hbar L_{x}$. Therefore

$$
\left[\hat{L}_{z} \hat{L}_{x}, \hat{L}_{y}\right]=L_{z}\left[L_{x}, L_{y}\right]+\left[L_{z}, L_{y}\right] L_{x}=i \hbar\left(L_{z}^{2}-L_{x}^{2}\right)
$$

vii) An unperturbed system has only two eigenfunctions $|1\rangle=x e^{-x^{2}}$ and $|2\rangle=x^{3} e^{-x^{2}}$ which are non-degenerate. Suppose there is a perturbing potential of the form $\alpha x$. Argue (i.e. do not calculate!) why this perturbation does not lead to a first or second order correction to either the energy or the wavefunctions. (3 pts)

Answer: All the integrals involved in the perturbation calculations are vanishing because we are integrating an odd function on a even domain. So the perturbation does not change the energetic levels.

## Problem 2 (16 pts)

Consider the wavefunction of the $|211\rangle$ level of the hydrogen atom

$$
\phi_{211}=\frac{1}{\sqrt{3}\left(2 a_{0}\right)^{3 / 2}} \frac{r}{a_{0}} \mathrm{e}^{-r / 2 a_{0}} Y_{1}^{1},
$$

with $Y_{1}^{1}=-\frac{1}{2} \sqrt{\frac{3}{2 \pi}} \sin (\theta) \mathrm{e}^{i \phi}$.
i) Show that $\phi_{211}$ is normalized. Hint: $\int_{0}^{\infty} d r r^{m} \mathrm{e}^{-x r}=m!x^{-(m+1)}$. ( 6 pts )

Answer:

$$
\begin{gathered}
\iiint \phi_{211}^{*} \phi_{211} d V=\int_{0}^{\infty} \int_{0}^{\pi} \int_{0}^{2 \pi} \frac{1}{3\left(2 a_{0}\right)^{3}} \frac{r^{2}}{a_{0}^{2}} \mathrm{e}^{-r / a_{0}} Y_{1}^{1 *} Y_{1}^{1} r^{2} \sin (\theta) d r d \theta d \phi= \\
\frac{1}{24} \int_{0}^{\infty} \rho^{4} e^{\rho} d \rho \times \frac{3}{8 \pi} \int_{0}^{\pi} \int_{0}^{2 \pi} \sin ^{3}(\theta) d \theta d \phi
\end{gathered}
$$

The integral over $\rho=r / a_{0}$ using repeated partial integration gives 1 , the second gives $\frac{3}{4} \int_{0}^{\pi} \sin ^{3}(\theta) d \theta=-\frac{3}{4} \int_{0}^{\pi}\left(1-\cos ^{2}(\theta)\right) d \cos (\theta)=1$
ii) What are the possible outcomes of a measurement of $L_{x}$ on this state? With what probabilities will these values occur?
Hint: Diagonalize $L_{x}$ and rewrite $Y_{1}^{1}$ in the new basis. (7 pts)
Answer: Eigenvectors-eigenvalues combinations are: $1 / \sqrt{2}(1,0,-1)$ with $0 \hbar, 1 / 2(1,-\sqrt{2}, 1)$ with $-\hbar$ and $1 / 2(1, \sqrt{2}, 1)$ with $\hbar$. Initial state is $\psi=$ $(1,0,0)$, so with $c_{m}=\left\langle\psi \mid \phi_{m}\right\rangle$ with $\phi_{m}$ the previous three normalized basis states we get $P=1 / 2$ for $0 \hbar$ and $P=1 / 4$ for $\pm \hbar$.
iii) Would an optical transition from this state to the state $|321\rangle$ be allowed by a dipole transition? Why/why not? (3 pts)

Answer: Yes, since it adheres to the selection rules $\Delta L=1, \Delta m=0$.

## Problem 3 (15 pts)

A two-level system is described by the following hamiltonian

$$
\hat{H}=\frac{\hbar \omega}{2}(-|0\rangle\langle 0|+|1\rangle\langle 1|),
$$

where $|0\rangle$ and $|1\rangle$ are the orthogonal eigenstates associated with the two energy levels ( $E_{0}<$ $E_{1}$ ).
i) Write $\hat{H}$ in matrix representation. (1 pt)

Answer:

$$
\hat{H}=\frac{\hbar \omega}{2}\left(\begin{array}{cc}
-1 & 0 \\
0 & 1
\end{array}\right)
$$

ii) What are the eigenvalues of this Hamilton operator? (1 pt)

Answer: The matrix is already diagonalized, so the eigenvalues are on the diagonal, so $\pm \hbar \omega / 2$.

Let's now introduce the operator $\hat{a}$ and its self-adjoint $\hat{a}^{\dagger}$. They act on the two levels as follows:

$$
\hat{a}|0\rangle=0 \quad \hat{a}|1\rangle=|0\rangle \quad \hat{a}^{\dagger}|0\rangle=|1\rangle \quad \hat{a}^{\dagger}|1\rangle=0 .
$$

iii) Write the matrix of $\hat{a}$ and $\hat{a}^{\dagger}$ in the basis of the eigenvectors of the hamiltonian. (2 pts)

Answer:

$$
\hat{a}^{\dagger}=\left(\begin{array}{cc}
\langle 0| a^{\dagger}|0\rangle & \langle 0| a^{\dagger}|1\rangle \\
\langle 1| a^{\dagger}|0\rangle & \langle 1| a^{\dagger}|1\rangle
\end{array}\right)=\left(\begin{array}{cc}
0 & 0 \\
1 & 0
\end{array}\right) \quad \hat{a}=\left(\begin{array}{cc}
0 & 1 \\
0 & 0
\end{array}\right)
$$

iv) Prove that $\hat{a} \hat{a}^{\dagger}+\hat{a}^{\dagger} \hat{a}=\mathbb{1}$, where $\mathbb{1}$ is the identity matrix. ( 2 pts )

Answer:

$$
\hat{a} \hat{a}^{\dagger}+\hat{a}^{\dagger} \hat{a}=\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right)\left(\begin{array}{ll}
0 & 0 \\
1 & 0
\end{array}\right)+\left(\begin{array}{ll}
0 & 0 \\
1 & 0
\end{array}\right)\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right)=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right)=\mathbb{1}
$$

We define now the operator $\hat{N}=\hat{a}^{\dagger} \hat{a}$ (number operator) and the operator $\hat{Q}=\hat{a}^{\dagger}-\hat{a}$.
v) Give the matrix representation for $\hat{Q}$ and find the eigenvalues and eigenvectors. (3 pts)

Answer:

$$
\begin{array}{cl}
Q=\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right) & q_{1}=i \quad q_{2}=-i \\
\left|q_{1}\right\rangle=\frac{1}{\sqrt{2}}\binom{i}{1} & \left|q_{2}\right\rangle=\frac{1}{\sqrt{2}}\binom{-i}{1}
\end{array}
$$

vi) Prove that the operators $\hat{N}$ represents a quantity that is conserved in time, and the the operator $\hat{Q}$ represents a quantity that is not conserved in time. (2 pts)

$$
\text { Answer: Since }[N, H]=0 \text { but }[Q, H] \neq 0, N \text { is conserved but } Q \text { is not. }
$$

vii) Suppose that the system is in the eigenstate $\left|q_{1}\right\rangle$ at time $t=0$. Evaluate explicitly $\langle Q\rangle$ for times $t>0$. ( 4 pts )

Answer:

$$
\left|q_{1}(t)\right\rangle=e^{-i H t / \hbar}\left|q_{1}\right\rangle=\frac{e^{-i H t / \hbar}}{\sqrt{2}}\left(i\binom{1}{0}+\binom{0}{1}\right)=\frac{1}{\sqrt{2}}\left[i e^{i \omega t / 2}|0\rangle+e^{-i \omega t / 2}|1\rangle\right] .
$$

With this outcome, we have

$$
\langle Q\rangle=i \cos (\omega t)
$$

## Problem 4 (12 pts)

The density operator that describes a beam of electrons, in the basis where $S_{z}$ is diagonal, has these elements:

$$
\rho_{11}=1 / 3 \quad \rho_{12}=(1+i) / 3 .
$$

i) What are the values of $\rho_{22}$ and $\rho_{21}$ ?

Using this, give the matrix representation of the density operator. (5 pts)
Answer: Since $\operatorname{Tr} \rho=1$, we achieve that $\rho_{22}=2 / 3$. Then, since $\rho$ has to be hermitian, we find that $\rho_{21}=(1-i) / 3$. Thus,

$$
\rho \doteq\left(\begin{array}{cc}
1 / 3 & (1+i) / 3 \\
(1-i) / 3 & 2 / 3
\end{array}\right)
$$

ii) A state is a pure state if and only if $\rho^{2}=\rho$. Is $\rho$ describing a pure state? ( 2 pts )

Answer: Since with this density matrix we have $\rho=\rho^{2}$, the system is in a pure state.
iii) Consider the matrices of $S_{x}$ and $S_{y}$ you found in Problem 1 ii). Evaluate the expectation values $\left\langle S_{x}\right\rangle$ and $\left\langle S_{y}\right\rangle$. (5 pts)

Answer: A fundamental property of the density operator is that $\langle A\rangle=$ $\operatorname{Tr}(\rho A)$. Therefore

$$
\begin{gathered}
\left\langle S_{x}\right\rangle=\operatorname{Tr}\left(\rho S_{x}\right)=\operatorname{Tr}\left[\left(\begin{array}{cc}
1 / 3 & (1+i) / 3 \\
(1-i) / 3 & 2 / 3
\end{array}\right)\left(\begin{array}{cc}
0 & \hbar / 2 \\
\hbar / 2 & 0
\end{array}\right)\right]=\hbar / 3 \\
\left\langle S_{y}\right\rangle=\operatorname{Tr}\left(\rho S_{x}\right)=\operatorname{Tr}\left[\left(\begin{array}{cc}
1 / 3 & (1+i) / 3 \\
(1-i) / 3 & 2 / 3
\end{array}\right)\left(\begin{array}{cc}
0 & -i \hbar / 2 \\
i \hbar / 2 & 0
\end{array}\right)\right]=-\hbar / 3
\end{gathered}
$$

## Problem 5 (25 pts)

An electron and a hole in a semiconductor have spin $1 / 2$. The Hamilton operator for them is given by

$$
\hat{H}_{0}=2 \gamma \hat{S}^{e} \cdot \hat{S}^{h}
$$

The superscript $e$ and $h$ stand for electron and hole, $\hat{S}$ are spin operators, and $\gamma$ is the spin-spin interaction strength. The eigenstates of the system are then a non-degenerate singlet state, and a triply degenerate triplet state. The singlet state is given by

$$
\begin{equation*}
\frac{1}{\sqrt{2}}(|\uparrow \downarrow\rangle-|\downarrow \uparrow\rangle) \tag{1}
\end{equation*}
$$

with energy $-\Delta(\Delta>0)$ and the triplet states are given by
$|\downarrow \downarrow\rangle$

$$
\begin{equation*}
\frac{1}{\sqrt{2}}(|\uparrow \downarrow\rangle+|\downarrow \uparrow\rangle) \tag{2}
\end{equation*}
$$

$|\uparrow \uparrow\rangle$
all three with energy 0 . The first arrow in the "ket's" stand for the electron spin, the second for the hole spin.
We call the normalized basis of these four states $\left\{\phi_{i}\right\}$.
i) Is $\gamma$ negative or positive? Why? ( 2 pts )

Answer: $\gamma$ is positive since the singlet state has lowest energy (i.e. the electron and hole spins energetically prefer not to be parallel aligned.

Now consider this system in the presence of a magnetic field $B$ in the $z$ direction. This leads to a additional term in the Hamilton operator of the form: $\hat{V}=\mu_{b} B\left(g_{e} \hat{S}_{z}^{e}+g_{h} \hat{S}_{z}^{h}\right)$, where $\mu_{b}$ is the Bohr magneton, and $g_{e}, g_{h}$ are the $g$-factors for the electron and hole, respectively. .
ii) Write down the action of the additional term $\hat{V}$ on the basis wavefunctions $\left\{\phi_{i}\right\}$ of $\hat{H}_{0}$. A useful definition here is $\alpha=\mu_{b} B\left(g_{e}-g_{h}\right) / 2$ and $\beta=\mu_{b} B\left(g_{e}+g_{h}\right) / 2$ ( 5 pts )

$$
\begin{aligned}
& \text { Answer: } \mathrm{V}|S\rangle=\alpha\left|T_{0}\right\rangle, \mathrm{V}\left|T_{0}\right\rangle=\alpha|S\rangle, \mathrm{V}\left|T_{-1}\right\rangle=-\beta\left|T_{-1}\right\rangle \text { and } \mathrm{V}\left|T_{+1}\right\rangle= \\
& \beta\left|T_{+1}\right\rangle .
\end{aligned}
$$

iii) Give the matrix representation of the Hamilton operator $\hat{H}=\hat{H}_{0}+\hat{V}$ using $\left\{\phi_{i}\right\}$ as basis. (3 pts)

$$
\text { Answer: }\left(\begin{array}{cccc}
-\Delta & 0 & \alpha & 0 \\
0 & -\beta & 0 & 0 \\
\alpha & 0 & 0 & 0 \\
0 & 0 & 0 & \beta
\end{array}\right)
$$

iv) Show that the diagonalization of this matrix gives the four eigenvalues $\pm \beta$ and $\frac{-\Delta \pm \sqrt{\Delta^{2}+4 \alpha^{2}}}{2}$. ( 5 pts )

Answer: The two eigenvalues $\pm \beta$ can be read directly from the matrix. What is left is to find the eigenvalues of $\left(\begin{array}{cc}-\Delta & \alpha \\ \alpha & 0\end{array}\right)$ which yields the two other values $\lambda_{ \pm}$.
v) Give the corresponding four eigenvectors for the Hamilton operator $\hat{H}$. ( 6 pts )

Answer: Using the basis $|S\rangle,\left|T_{-1}\right\rangle,\left|T_{0}\right\rangle,\left|T_{+1}\right\rangle$, the eigenvectors-eigenvalues pairs are $(0,1,0,0)$ with $-\beta,(0,0,0,1)$ with $\beta,\left(1,0, \alpha / \lambda_{+}, 0\right)$ with $\lambda_{+}$and $\left(1,0, \alpha / \lambda_{-}, 0\right)$ with $\lambda_{-}$.
vi) Assume that $|\alpha| \ll|\Delta|$, give an approximation for the wavefunctions (valid to order $\alpha / \Delta$ ) and eigenvalues (valid to order $\alpha^{2} / \Delta$ ). (3 pts)

Answer: $(0,1,0,0)$ with $-\beta,(0,0,0,1)$ with $\beta,(1,0,-\alpha / \Delta, 0)$ with $-\Delta-\frac{\alpha^{2}}{\Delta}$ and $(\alpha / \Delta, 0,1,0)$ with $\frac{\alpha^{2}}{\Delta}$
vii) How do the four energy levels scale with the magnetic field? (1 pts)

Answer: Singlet and $m=0$ of the triplet scale quadratic in B, the other two linear in B.

